

# 10<sup>th</sup> ABINIT International Developer Workshop

May 31 - June 4, 2021

Are plane-waves basis packages able to predict the correct long range behavior of the local kinetic energy per particle?

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Camilo Espejo, Hugo Bohórquez and Aldo Romero





# Summary

1. Background and Motivation.
2. Two theorems involving  $n(\mathbf{r})$  and  $G(\mathbf{r})$  in finite systems.
3. Computing the asymptotic behavior of  $G(\mathbf{r})$  in **ABINIT** for atoms and molecules.
4. Conclusions and perspectives.

# Background and Motivation



nature  
materials

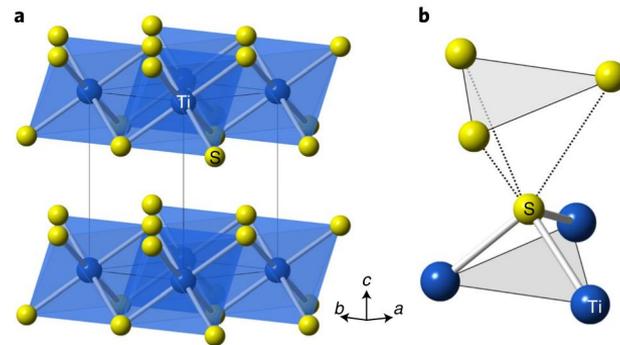
ARTICLES

<https://doi.org/10.1038/s41563-017-0012-2>

## X-ray electron density investigation of chemical bonding in van der Waals materials

Hidetaka Kasai<sup>1,2</sup>, Kasper Tolborg<sup>2</sup>, Mattia Sist<sup>2</sup>, Jiawei Zhang<sup>2</sup>, Venkatesha R. Hathwar<sup>1</sup>, Mette Ø. Filsø<sup>2</sup>, Simone Cenedese<sup>2</sup>, Kunihisa Sugimoto<sup>3</sup>, Jacob Overgaard<sup>2</sup>, Eiji Nishibori<sup>1</sup> and Bo B. Iversen<sup>2\*</sup>

Nature Materials 17:249 (2018)



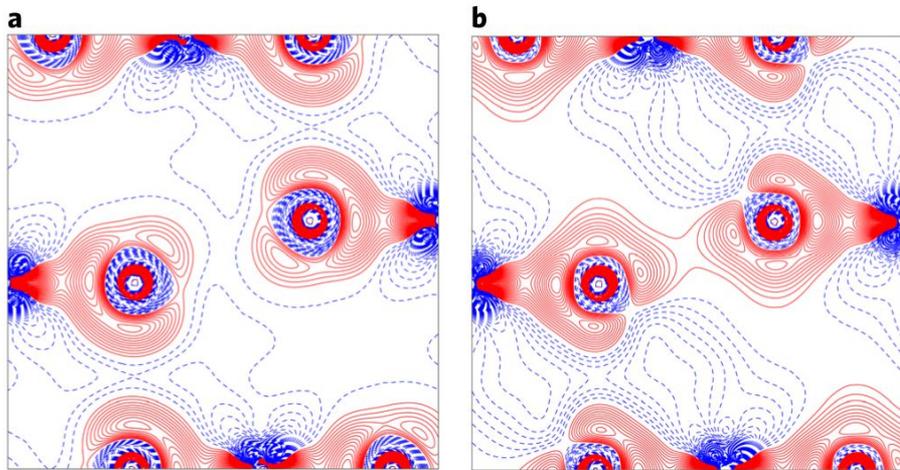
**Fig. 1 | Structure of 1T-TiS<sub>2</sub>.** **a**, Ti atom is at the edge of a unit cell (Space group:  $P\bar{3}m1$ ,  $a = 3.3976(3)$  and  $c = 5.6654(3)$  Å) forming a TiS<sub>6</sub> octahedron. **b**, S atom and its adjacent atoms in the direction of the  $c$ -axis.

# Background and Motivation

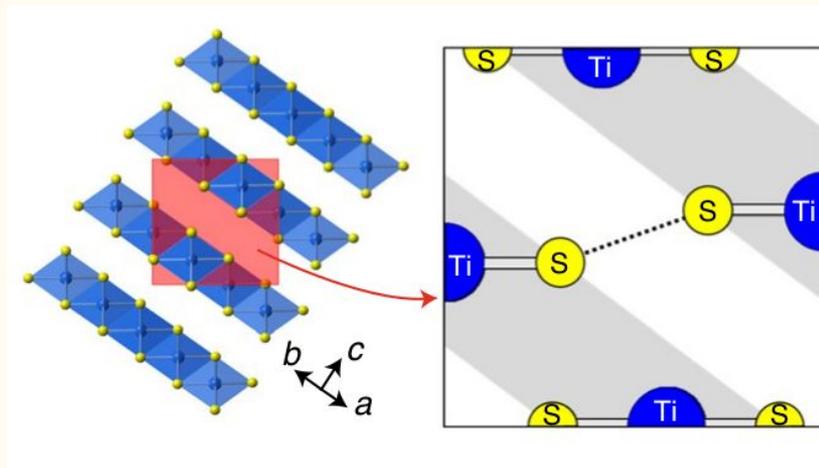


Experiment

Theory(SCAN+rVV10)



**interaction. a,b**, Static deformation map on the Ti-S-S-Ti plane (110 plane) from multipole modelling of the theoretical (SCAN+rVV10) (a) and experimental (b) structure factors. The contour interval is  $0.01 e \text{ \AA}^{-3}$ , with positive and negative contours drawn as solid red and dotted blue lines, respectively. **c,d**, Negative Laplacian map obtained



Even though both equilibrium interlayer distance and binding energy follow the experimental trend, there are important differences in the ground state electronic density.

# Background and motivation



**Table 2 | Atomic properties of Ti and S**

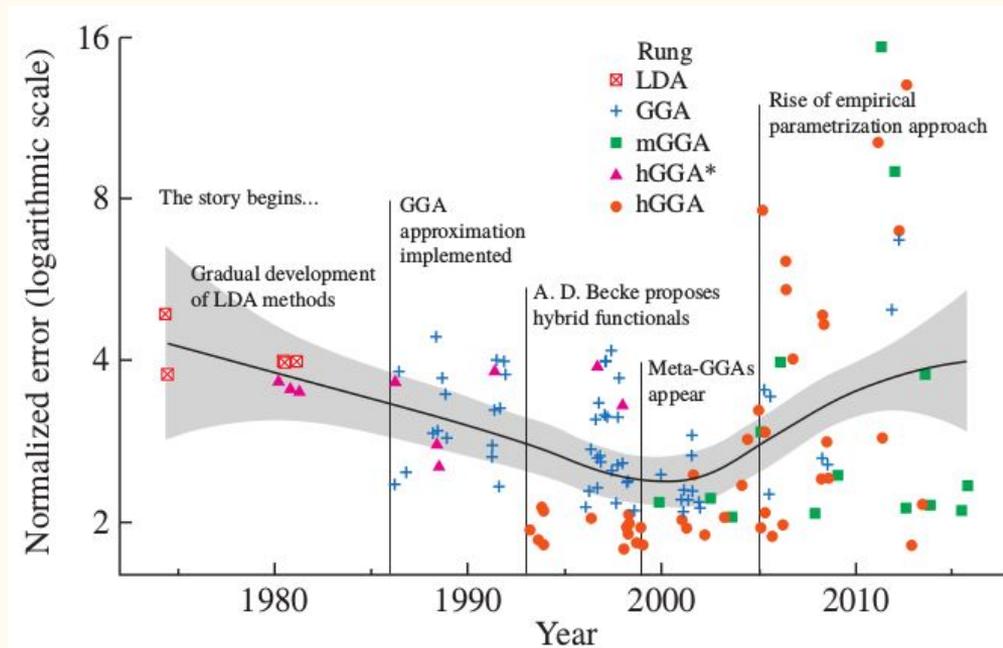
Variable	Method	Ti	S
Q	Experiment	1.63	-0.82
	Theory (SCAN+rVV10)	1.58	-0.80
V (Å <sup>3</sup> )	Experiment	9.89	23.47
	Theory (SCAN+rVV10)	9.64	23.63
$\mu$ (eÅ)	Experiment	0.00	0.03
	Theory (SCAN+rVV10)	0.00	0.34
L (a.u.)	Experiment	$-7.47 \times 10^{-3}$	$-4.34 \times 10^{-4}$
	Theory (SCAN+rVV10)	$-1.05 \times 10^{-2}$	$-1.36 \times 10^{-4}$

$Q$ ,  $V$  and  $\mu$  are the atomic charge, the atomic basin volume and the magnitude of atomic dipole moment, respectively.  $L$  is the atomic Lagrangian  $(-(1/4) \int_{\Omega} \nabla^2 \rho \, d\tau)$ , which is a measure of the accuracy of the integration.

Nature Materials 17:249 (2018)

# Houston, we have a problem!

Medvedev, *et. al. Science* **355**, 49-52(2017)



REPORT

THEORETICAL CHEMISTRY

## Density functional theory is straying from the path toward the exact functional

Michael G. Medvedev,<sup>1,2,3\*†</sup> Ivan S. Bushmarinov,<sup>1\*†</sup> Jianwei Sun,<sup>4‡</sup> John P. Perdew,<sup>4,5†</sup> Konstantin A. Lyssenko<sup>1†</sup>

Mendelev Commun., 2018, 28, 225–235

“The reproduction of the exact energy is not a characteristic of the exact functional unless the input electronic density is exact as well.”

# Houston, we have a problem!

Medvedev, *et. al. Science* **355**, 49-52(2017)

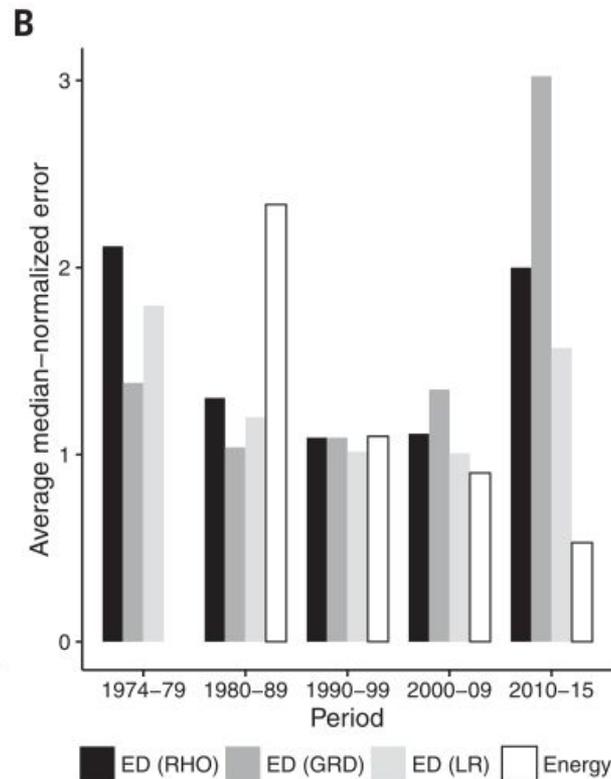
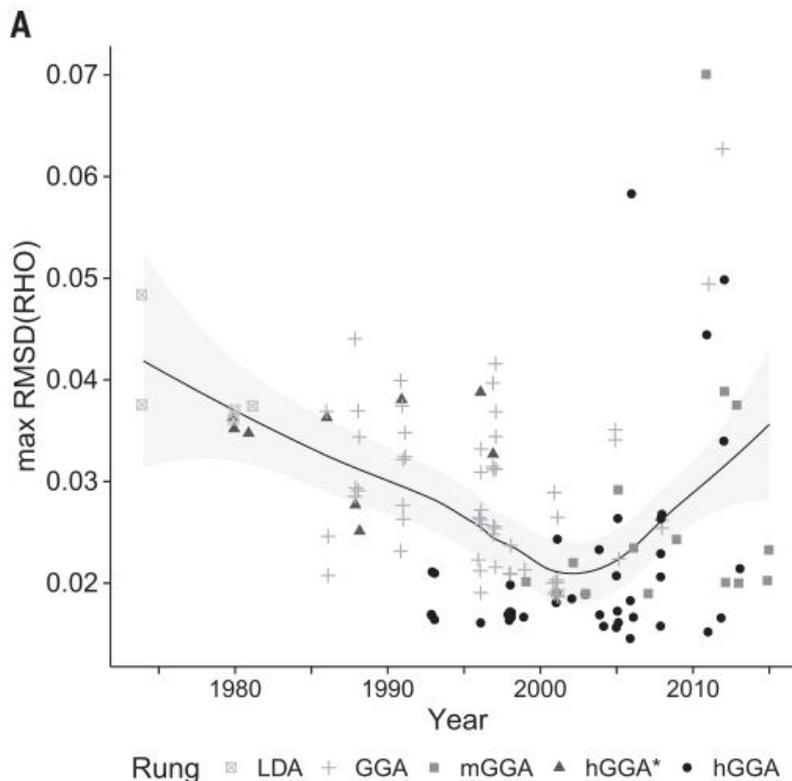


REPORT

THEORETICAL

## Density functional theory: the evolution of the exchange-correlation functional

Michael G. Medvedev  
John P. Perdew



# Measuring the accuracy of the electronic density



Two main groups of methods:

- Metrics based on integrals of the density:

Density moments, Total kinetic energy.

Problems can arise due to possible cancellation of errors!

- Metrics based on point to point measurements. (Local values)

# Measuring the accuracy of the electronic density



Some metrics based on point to point comparisons:

$$\text{RMSD} = \sqrt{\frac{1}{N} \sum_i^N \{ \pi r_i^2 [P_{\text{approx}}(r_i) - P_{\text{ref}}(r_i)] \}^2},$$

Root mean-square deviation

$$\text{ISD} = \iiint [P_{\text{approx}}(r) - P_{\text{ref}}(r)]^2 d^3r,$$

Integral squared deviation.

$$\text{IAD} = \iiint |P_{\text{approx}}(r) - P_{\text{ref}}(r)| d^3r.$$

Integral absolute deviation.

$P$  can be any local descriptor such as the density, norm of the gradient, the laplacian, etc...

2. Two theorems for  $n(\mathbf{r})$  and  $G(\mathbf{r})$  in finite systems.



Kinetic energy density:

$$\tau_{ks} = \frac{1}{2} \sum_i^{\text{occup}} f_i |\nabla \phi_i|^2$$

Kinetic energy per particle:

$$G(\mathbf{r}) = \frac{\tau_{ks}(\mathbf{r})}{n(\mathbf{r})}$$

## 2. Two theorems for the single particle density in finite systems.



From Kato's theorem (Kato cusp condition):  $Z_k = -\frac{a_0}{2n(\mathbf{r})} \frac{dn(\mathbf{r})}{dr} \Big|_{r \rightarrow R_k}$

$$\lim_{r \rightarrow R_k} G(\mathbf{r}) = \frac{Z_k^2}{2}$$

*Communications on Pure and Applied Mathematics.* **10** (2): 151–177 (1957)

## 2. Two theorems for the single particle density in finite systems.



From Hoffmann-Ostenhof theorems for the asymptotic behavior of the multi-electronic wave function of finite coulombic systems:

$$\lim_{r \rightarrow \infty} G(\mathbf{r}) = I$$

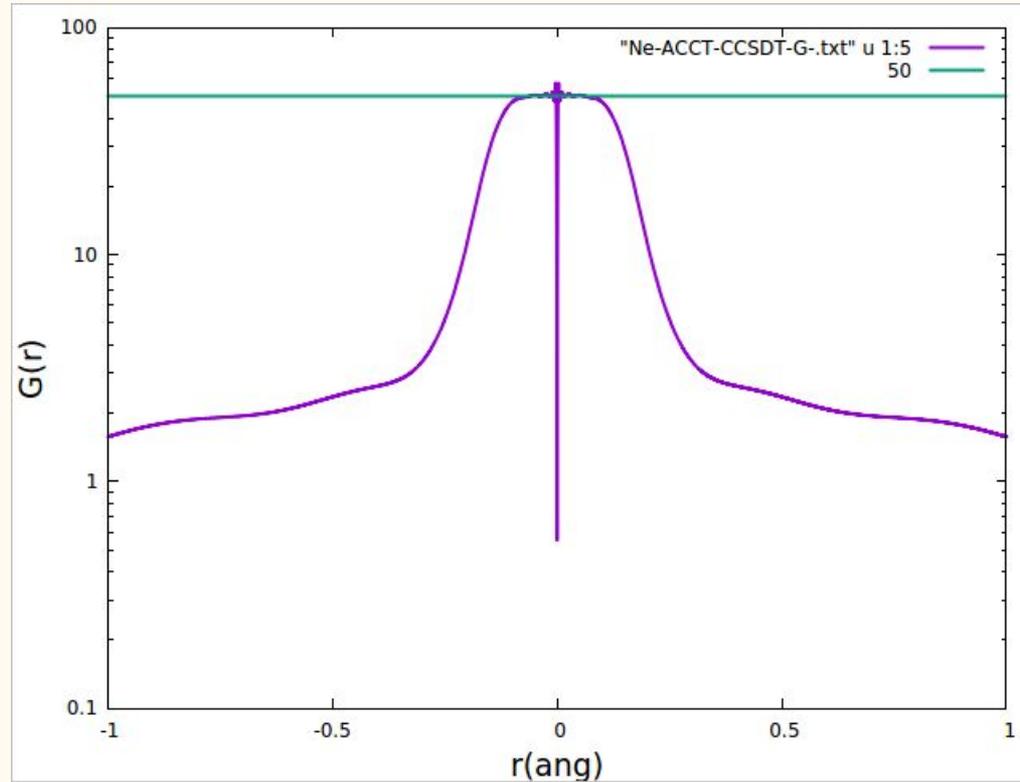
*J. Phys.B: Atom. Molec. Phys.*, **11:17** (1978), *Chem. Phys. Lett.* **258:30** (1996)

### 3. Computing the asymptotic behavior of $G(\mathbf{r})$ in ABINIT for atoms and molecules.



First some reference calculations: Ne atom CCSDT (ACCT) in GAMESS

$$Z_{\text{Ne}} = 10$$

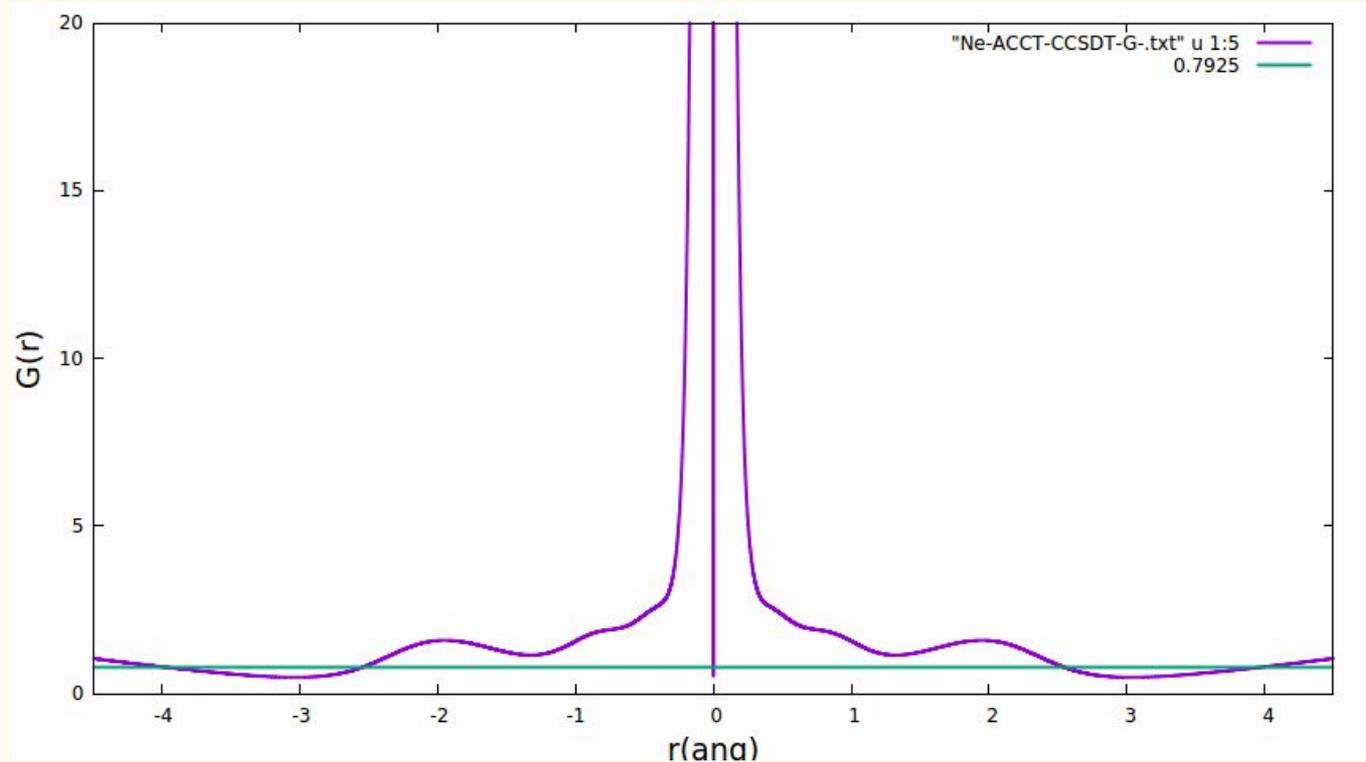


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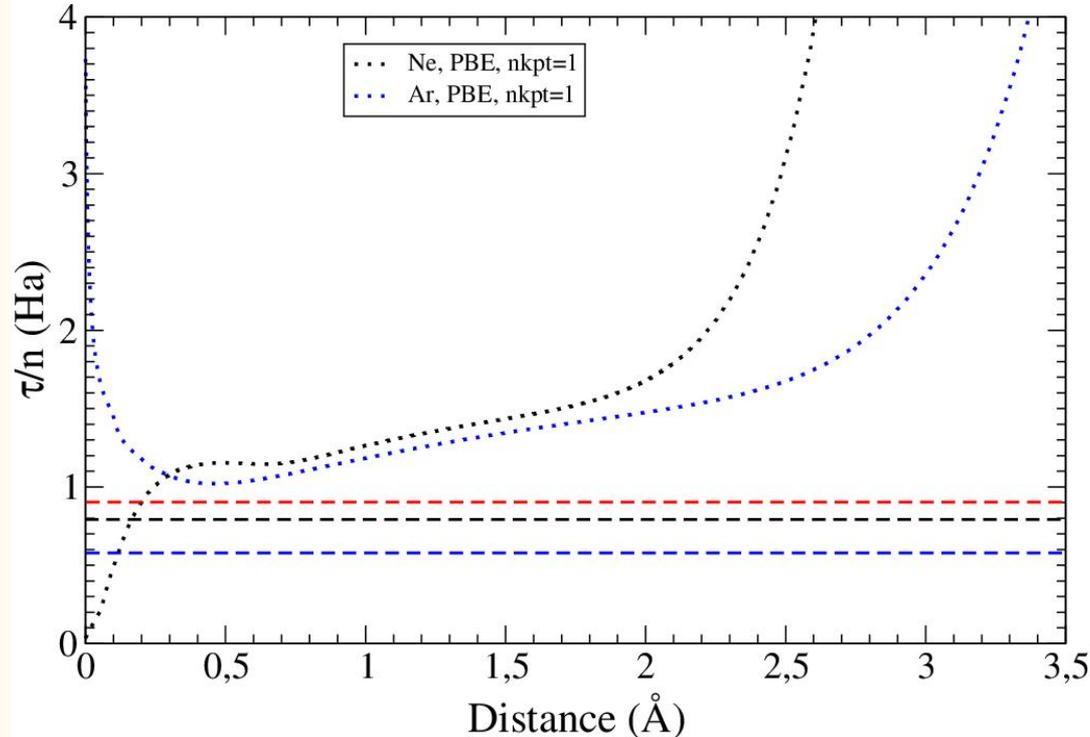
$I=0.7925$  Ha



### 3. Computing the asymptotic behavior of $G(\mathbf{r})$ in ABINIT for atoms and molecules.



Noble gases

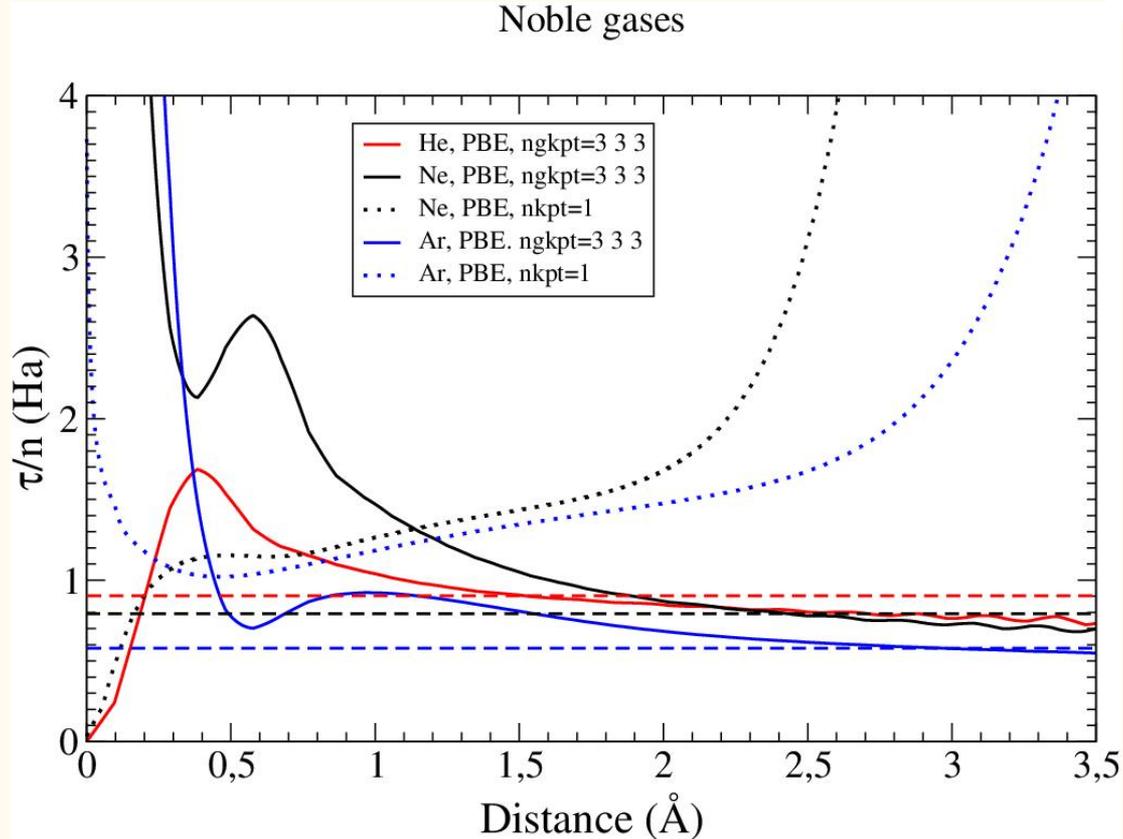


ecut=80 Ha  
acell=3\*15 A

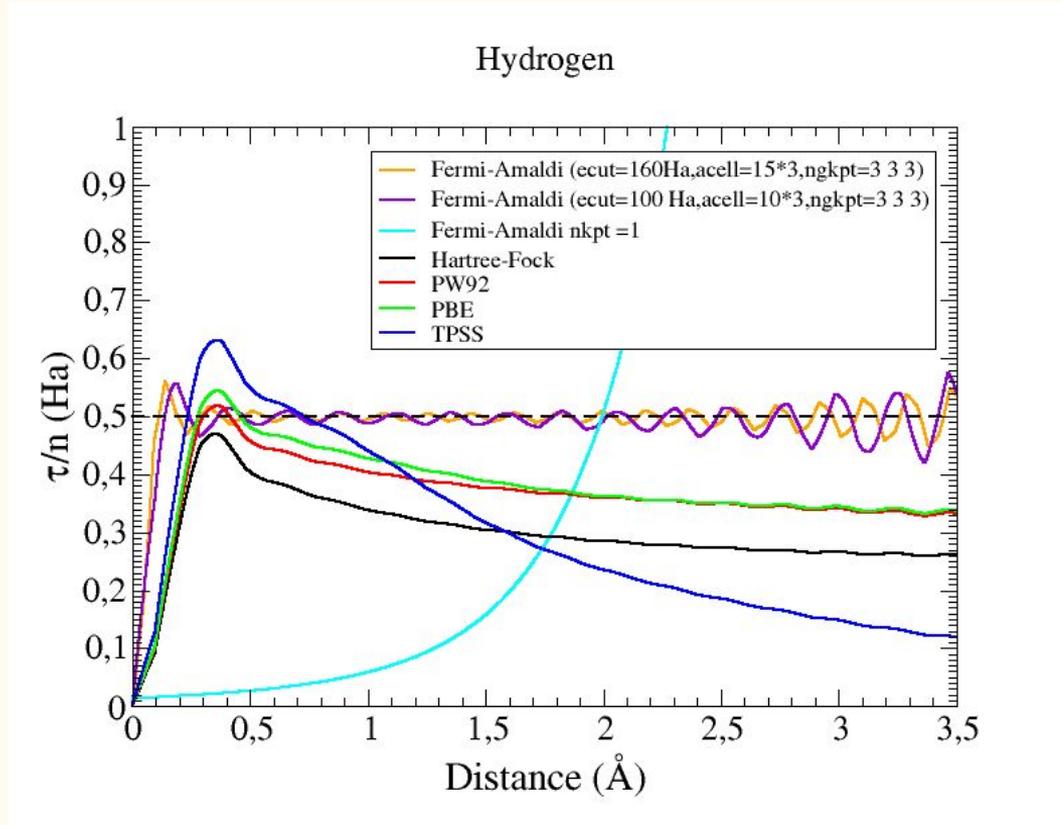
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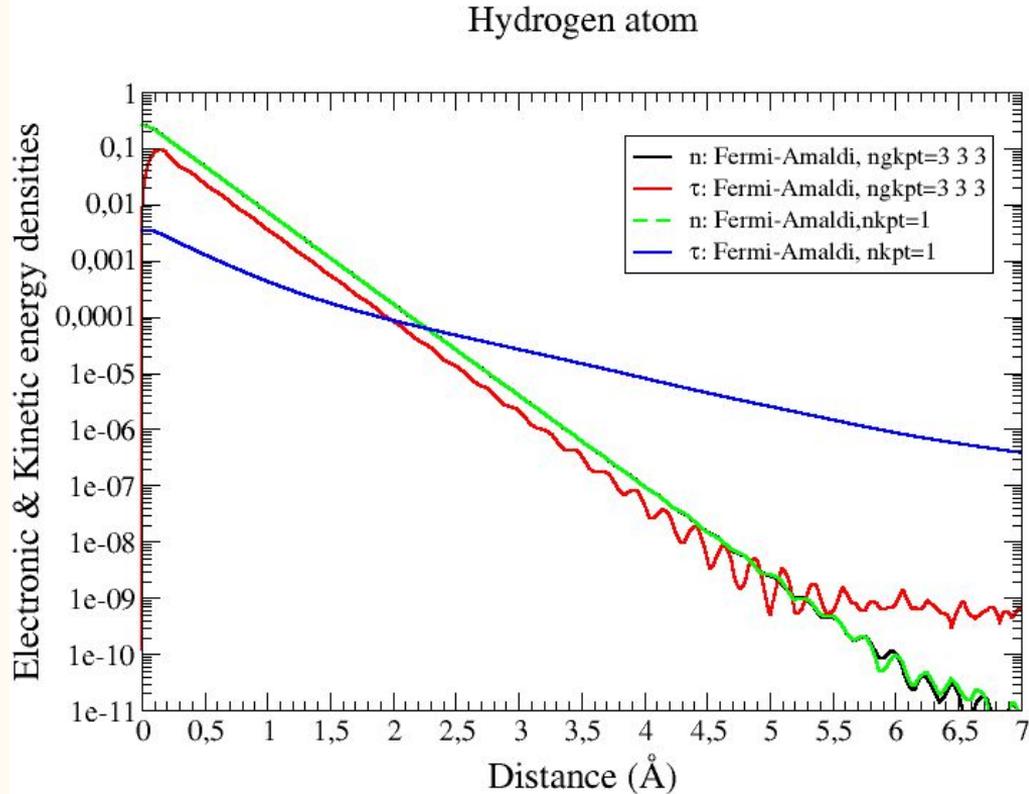
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acell=3\*15 A



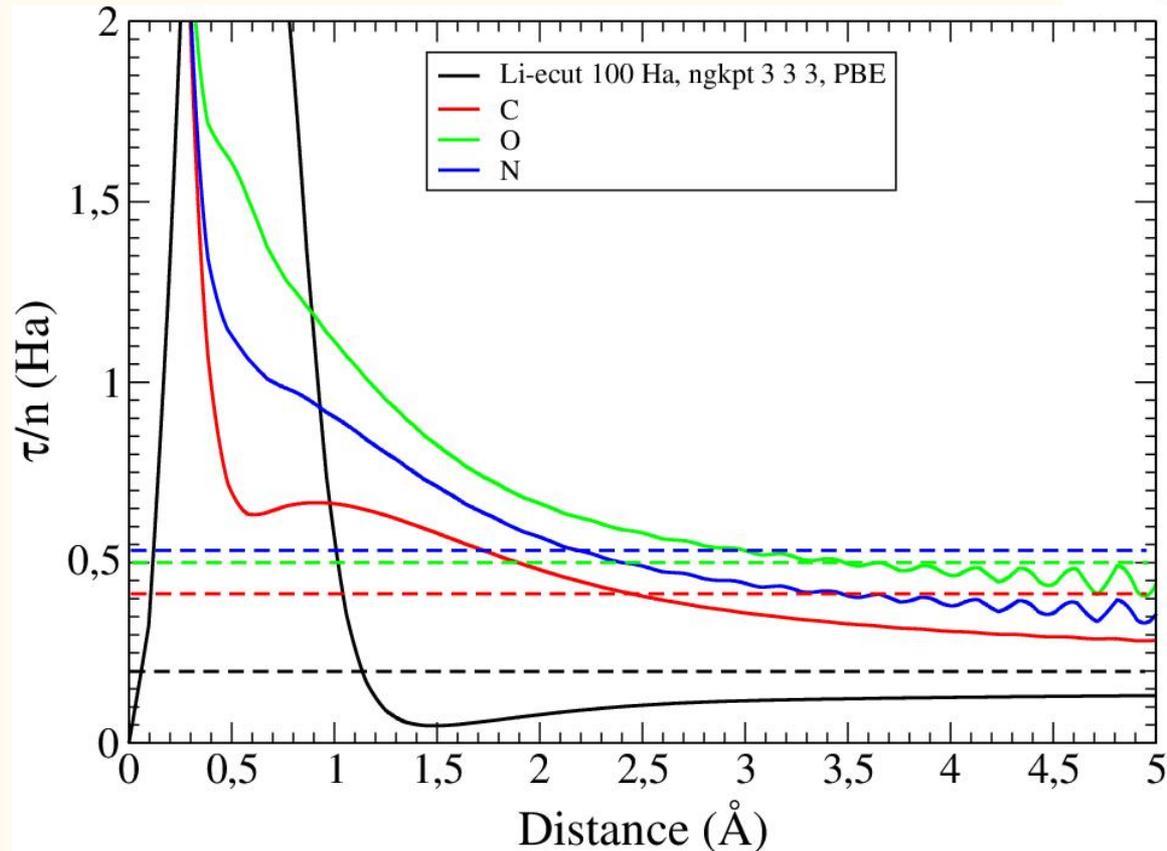
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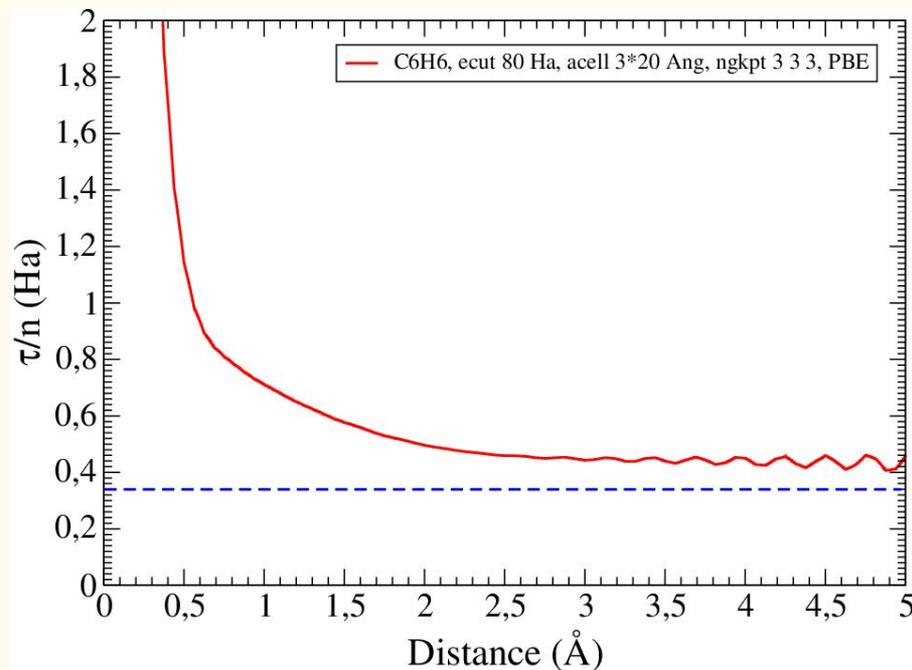
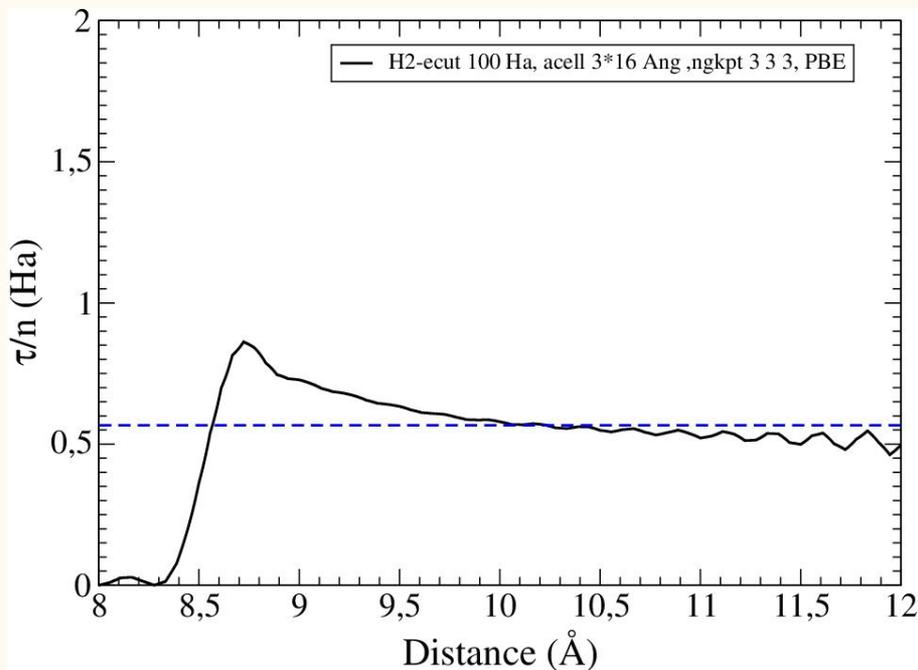
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# 4. Conclusions and perspectives.

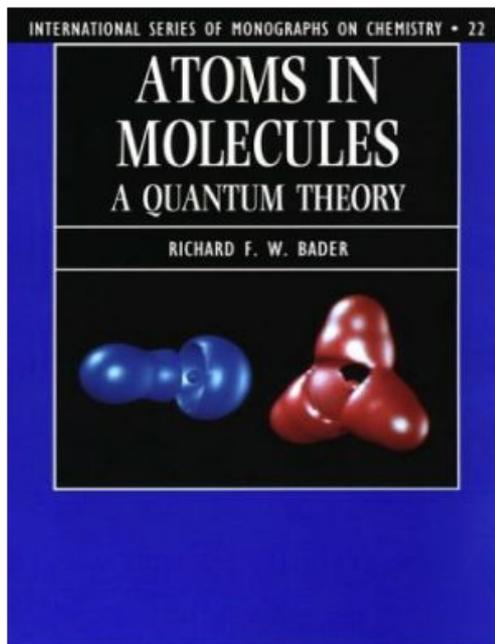
- It is possible to obtain a reasonable behavior of  $G(\mathbf{r})$  within the ncpp + plane-waves framework for finite systems if a fine enough  $k$  sampling is used.
- Since  $\tau$  in real space is an input for meta-GGA functionals it is important to add  $k$  points convergence studies when dealing with finite systems (i.e. adsorption of molecules on surfaces).
- $G(\mathbf{r})$  can play an important role in the benchmarking of vdW self-consistent functionals from the density point of view.
- What could be the effect of the Coulomb truncation methods (from the last talk) on the asymptotic behavior of  $G$ ?
- What is the meaning of the limit of  $G$  in extended systems (2D materials, surfaces)??



Thank you for your attention!

# Appendix

## QTAIM Basics (Bader and Co)



# QTAIM Basics (Bader and Co)

$$\nabla\rho = \mathbf{i} \frac{d\rho}{dx} + \mathbf{j} \frac{d\rho}{dy} + \mathbf{k} \frac{d\rho}{dz} \rightarrow \begin{cases} = \vec{\mathbf{0}} & \text{(At critical points and} \\ & \text{at } \infty) \\ \text{Generally } \neq \vec{\mathbf{0}} & \text{(At all other points)} \end{cases}$$

$$\mathbf{A}(\mathbf{r}_c) = \begin{pmatrix} \frac{\partial^2 \rho}{\partial x^2} & \frac{\partial^2 \rho}{\partial x \partial y} & \frac{\partial^2 \rho}{\partial x \partial z} \\ \frac{\partial^2 \rho}{\partial y \partial x} & \frac{\partial^2 \rho}{\partial y^2} & \frac{\partial^2 \rho}{\partial y \partial z} \\ \frac{\partial^2 \rho}{\partial z \partial x} & \frac{\partial^2 \rho}{\partial z \partial y} & \frac{\partial^2 \rho}{\partial z^2} \end{pmatrix}_{\mathbf{r}=\mathbf{r}_c}$$

# QTAIM Basics (Bader and Co)

$$\Lambda = \begin{pmatrix} \frac{\partial^2 \rho}{\partial x'^2} & 0 & 0 \\ 0 & \frac{\partial^2 \rho}{\partial y'^2} & 0 \\ 0 & 0 & \frac{\partial^2 \rho}{\partial z'^2} \end{pmatrix}_{\mathbf{r}'=\mathbf{r}_c} = \begin{pmatrix} \lambda_1 & 0 & 0 \\ 0 & \lambda_2 & 0 \\ 0 & 0 & \lambda_3 \end{pmatrix},$$

$$\nabla^2 \rho(\mathbf{r}) = \nabla \cdot \nabla \rho(\mathbf{r}) = \underbrace{\frac{\partial^2 \rho(\mathbf{r})}{\partial x^2}}_{\lambda_1} + \underbrace{\frac{\partial^2 \rho(\mathbf{r})}{\partial y^2}}_{\lambda_2} + \underbrace{\frac{\partial^2 \rho(\mathbf{r})}{\partial z^2}}_{\lambda_3}$$

# QTAIM Basics (Bader and Co)

There are four types of stable critical points having three non-zero eigenvalues:

- $(3, -3)$  Three negative curvatures:  $\rho$  is a local maximum.
- $(3, -1)$  Two negative curvatures:  $\rho$  is a maximum in the plane defined by the corresponding eigenvectors but is a minimum along the third axis which is perpendicular to this plane.
- $(3, +1)$  Two positive curvatures:  $\rho$  is a minimum in the plane defined by the corresponding eigenvectors and a maximum along the third axis which is perpendicular to this plane.
- $(3, +3)$  Three curvatures are positive:  $\rho$  is a local minimum.

# QTAIM Basics (Bader and Co)

(3,-3)	All curvatures -ve, a local maximum
(3,-1)	Two curvatures are -ve and one is +ve $\rho$ is a maximum in a plane and a minimum perpendicular to this plane - a <b>bond cp</b>
(3,+1)	Two curvatures are +ve and one is -ve $\rho$ is a minimum in a plane and a maximum perpendicular to this plane - a <b>ring cp</b>
3,+3)	All curvatures +ve, a local minimum - a <b>cage cp</b>

# QTAIM Basics (Bader and Co)

Poincaré-Hopf relationships:

$$n_{\text{NCP}} - n_{\text{BCP}} + n_{\text{RCP}} - n_{\text{CCP}} = \begin{cases} 1 & \text{(Isolated molecules)} \\ 0 & \text{(Infinite crystals)} \end{cases}$$

# QTAIM Basics (Bader and Co)

